

Derivation of Brownian Motion from the Random Walk

Michael Hauser

This article gives a short derivation of Brownian motion from the random walk in an easy to follow way. First define the step random variable Z_k , where we take a step to the right with probability p and we take a step to the left with probability $1 - p$:

$$Z_k := \begin{cases} +\Delta x & \text{with probability } p \\ -x & \text{with probability } q = 1 - p \end{cases} \quad (1)$$

Then the total distance traveled is the sum:

$$S_n := \sum_{k=1}^n Z_k \quad (2)$$

To find the mean and variance of S_n , use the characteristic function, where $E[\cdot]$ is the expectation:

$$\begin{aligned} \varphi_{S_n}(u) &:= E[e^{iS_n u}] = E[e^{i\sum_{k=1}^n Z_k u}] = \\ E\left[\prod_{k=1}^n e^{iZ_k u}\right] &= \prod_{k=1}^n E[e^{iZ_k u}] = E[e^{iZ_k u}]^n = \\ &= (pe^{i\Delta x u} + qe^{-i\Delta x u})^n \end{aligned} \quad (3)$$

Let there be a total of $n = \frac{t}{\Delta t}$ steps over time t and time-step Δt , then:

$$\varphi_{S_n}(u) = (pe^{i\Delta x u} + qe^{-i\Delta x u})^{t/\Delta t} \quad (4)$$

Remark. (Moment Generating Functions)
 $\varphi_x(k) = E[e^{ikx}] \Rightarrow \frac{d^m \varphi_x}{dk^m}(k) = \left(\frac{d}{dk}\right)^m E[e^{ikx}] = E[(ix)^m e^{ikx}]$. This means that at $k = 0$ we have: $\frac{d^m \varphi_x}{dk^m}(k = 0) = i^m E[x^m]$.

From the remark, we have the result:

$$E[x^m] = \frac{1}{i^m} \frac{d^m \varphi_x}{dk^m}(0) \forall m \in \mathbb{N} \quad (5)$$

With this, and letting $p = q = 1/2$:

$$E[(S_n)^m] = \frac{1}{i^m} \frac{d^m \varphi_{S_n}}{du^m}(u = 0) \forall m \in \mathbb{N} \quad (6)$$

For $m = 1$ (i.e. the mean):

$$E[S_n] = \frac{1}{i} \frac{d\varphi_{S_n}}{du}(0) = \frac{1}{i} \left(\frac{1}{2}\right)^{t/\Delta t} \xrightarrow{\Delta t \rightarrow 0} 0 \quad (7)$$

For $m = 2$ (i.e. the variance):

$$\begin{aligned} E[(S_n)^2] &= \frac{1}{i^2} \frac{d^2 \varphi_{S_n}}{du^2}(0) = \dots \text{lots of algebra} \dots = \\ &= t \frac{\Delta x^2}{\Delta t} \xrightarrow{\Delta x, \Delta t \rightarrow 0} t\sigma^2 \end{aligned} \quad (8)$$

Thus we have, with the limiting definition (an example of Einstein's genius) $\frac{\Delta x^2}{\Delta t} \xrightarrow{\Delta x, \Delta t \rightarrow 0} \sigma^2$. Using these results, we substitute this limiting definition into Equation 4 (along with $p = q = 1/2$) to generate the density function in the continuous limit:

$$\begin{aligned} \varphi_{S_n}(u) &= \left(\frac{1}{2} (e^{i\Delta x u} + e^{-i\Delta x u})\right)^{t/\Delta t} \xrightarrow{\Delta x, \Delta t \rightarrow 0} \\ \varphi_{x_t}(u) &= \left(\frac{1}{2} (e^{iu\sigma\sqrt{\Delta t}} + e^{-iu\sigma\sqrt{\Delta t}})\right)^{t/\Delta t} = \\ &= \cos(u\sigma\sqrt{\Delta t})^{t/\Delta t} \end{aligned} \quad (9)$$

Note that the name of the *discrete* stochastic variable is S_n , whereas after the limit we have the *continuous* stochastic variable x_t .

Now do two Taylor expansions and ignore the higher order terms. The first Taylor expansion:

$$\begin{aligned} \varphi_{x_t}(u) &= \cos\left(u\sigma\sqrt{\Delta t}\right)^{t/\Delta t} = \\ &\left(1 - \frac{1}{2}\left(u\sigma\sqrt{\Delta t}\right)^2 + \mathcal{O}\left(\sqrt{\Delta t}^3\right)\right)^{t/\Delta t} \end{aligned} \quad (10)$$

The second Taylor expansion:

$$\begin{aligned} \log \varphi_{x_t}(u) &= \log\left(1 - \frac{1}{2}u^2\sigma^2\Delta t\right)^{t/\Delta t} = \\ &\frac{t}{\Delta t} \log\left(1 - \frac{1}{2}u^2\sigma^2\Delta t\right) = \\ &\frac{t}{\Delta t} \left(-\frac{1}{2}u^2\sigma^2\Delta t + \mathcal{O}(\Delta t)^2\right) \end{aligned} \quad (11)$$

We can finally rearrange this final term as follows:

$$\begin{aligned} \log \varphi_{x_t}(u) &= -\frac{1}{2}u^2\sigma^2 t \iff \\ \varphi_{x_t}(u) &= e^{-\frac{1}{2}u^2\sigma^2 t} \iff \\ &x_t \sim N(0, \sigma^2 t) \end{aligned} \quad (12)$$

Conclusion: We have our result, the discrete random walk turns into Brownian motion in the continuous limit.